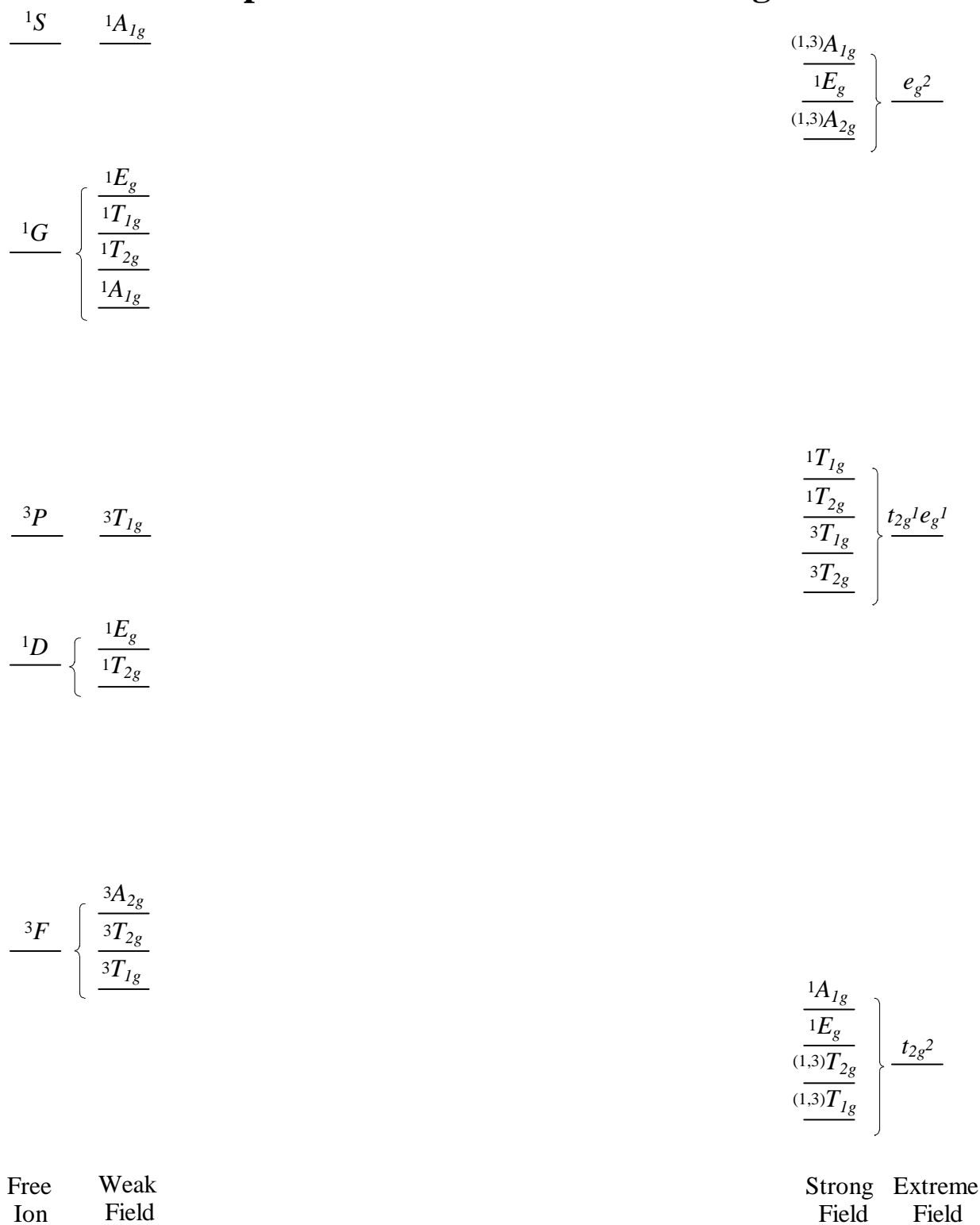


## Template for the $d^2$ Correlation Diagram



## Making the Correlations for Triplet States

### Correlation for the ${}^3A_{2g}$ Term from ${}^3F$

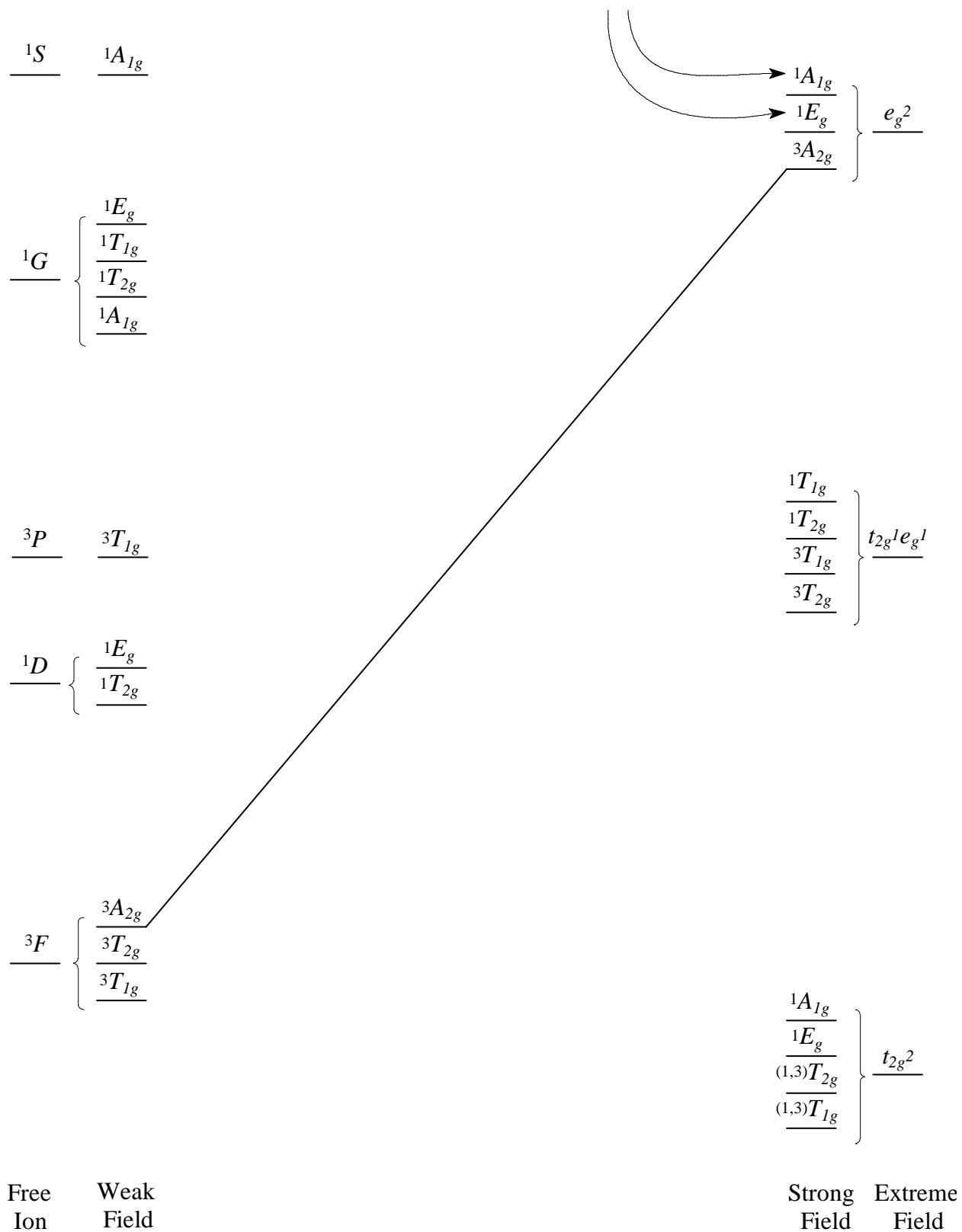
- We begin by examining the split terms from the  ${}^3F$  free-ion term on the left of the diagram.
  - The  ${}^3A_{2g}$  term is unique among the weak-field terms, and it must correlate with the unique  $A_{2g}$  term from  $e_g^2$  on the strong-field side.
    - ☞ The strong-field  $A_{2g}$  term must also must be a triplet.
    - ☞ It now becomes evident that our choice of spin multiplicities for the terms from  $e_g^2$  must be

$${}^1A_{1g} + {}^3A_{2g} + {}^1E_g$$

- We can draw the correlation line for the  ${}^3A_{2g}$  term and remove the ambiguous spin state notation for the  $A_{1g}$  and  $E_g$  terms for  $e_g^2$  on the strong-field side of the diagram.

# Correlation for $^3A_{2g}$ Added

These spin states can now be established.



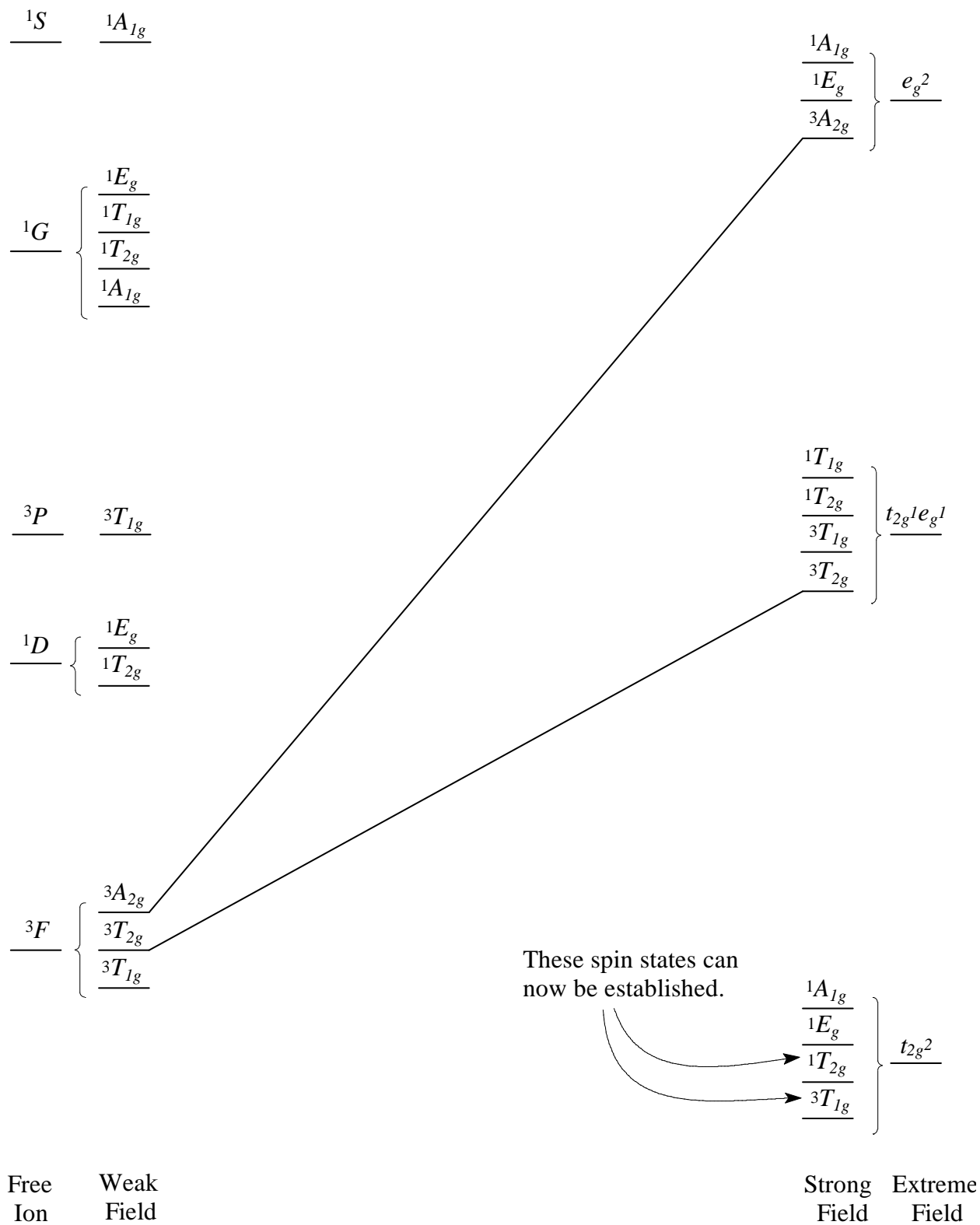
## Making the Correlations for Triplet States

### Correlation for the ${}^3T_{2g}$ Term from ${}^3F$

- We turn now to the  ${}^3T_{2g}$  term arising from  ${}^3F$ .
  - This is the only triplet  $T_{2g}$  term on the weak-field side (the two other  $T_{2g}$  terms are singlets).
    - ☞ There can be only one such term on the strong-field side.
  - We have already identified a  ${}^3T_{2g}$  term from  $t_{2g}{}^1e_g{}^1$  on the strong-field side, which must correlate with the  ${}^3T_{2g}$  term from  ${}^3F$ .
    - ☞ This means that the  $T_{2g}$  term from  $t_{2g}{}^2$  on the strong field side must be a singlet.
    - ☞ From this it follows that the spin multiplicities for terms from  $t_{2g}{}^2$  must be

$${}^1A_{1g} + {}^1E_g + {}^3T_{1g} + {}^1T_{2g}$$

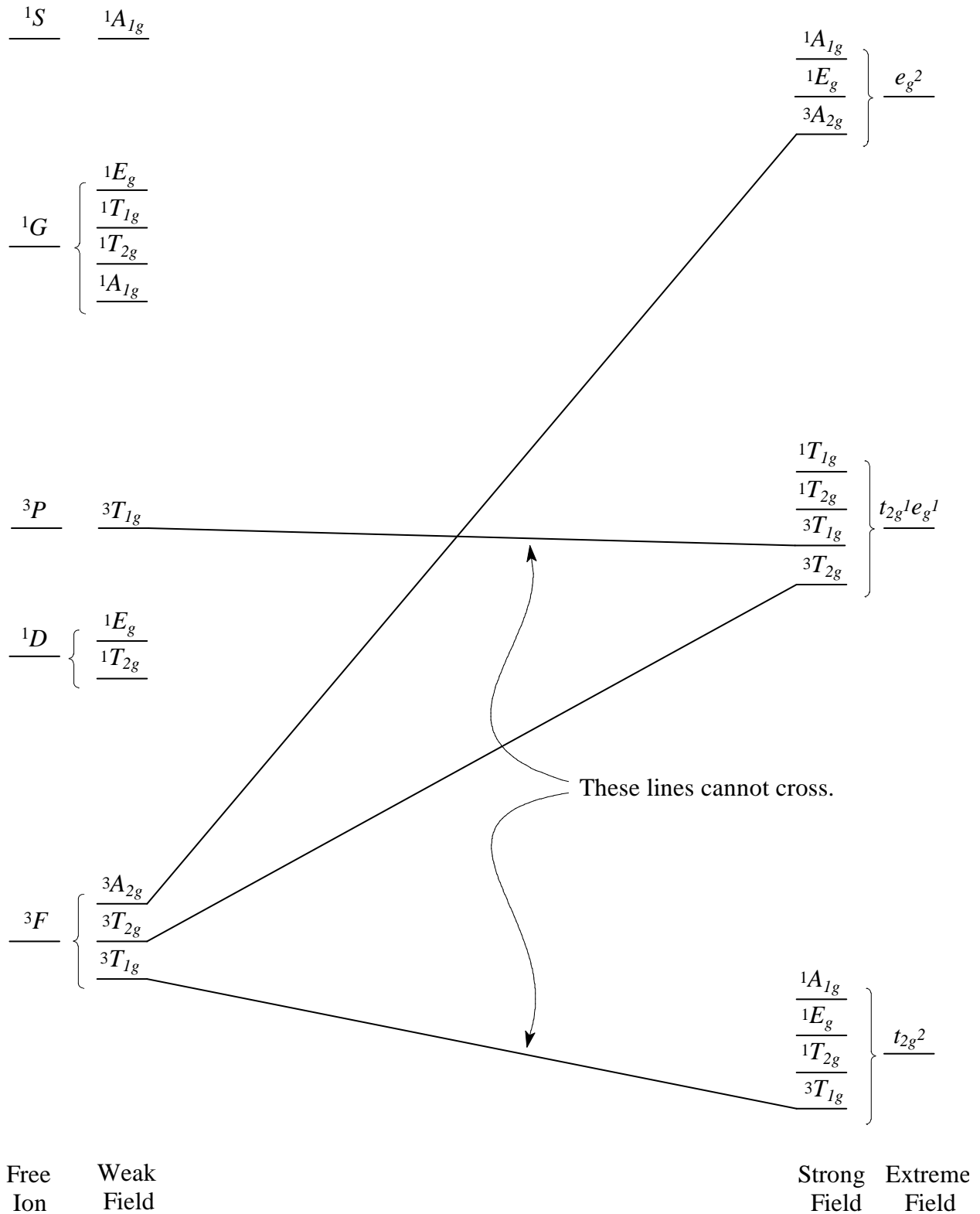
# Correlation for ${}^3T_{2g}$ Added



## Making the Correlations for Triplet States Correlation for the ${}^3T_{1g}$ Terms from ${}^3F$ and ${}^3P$

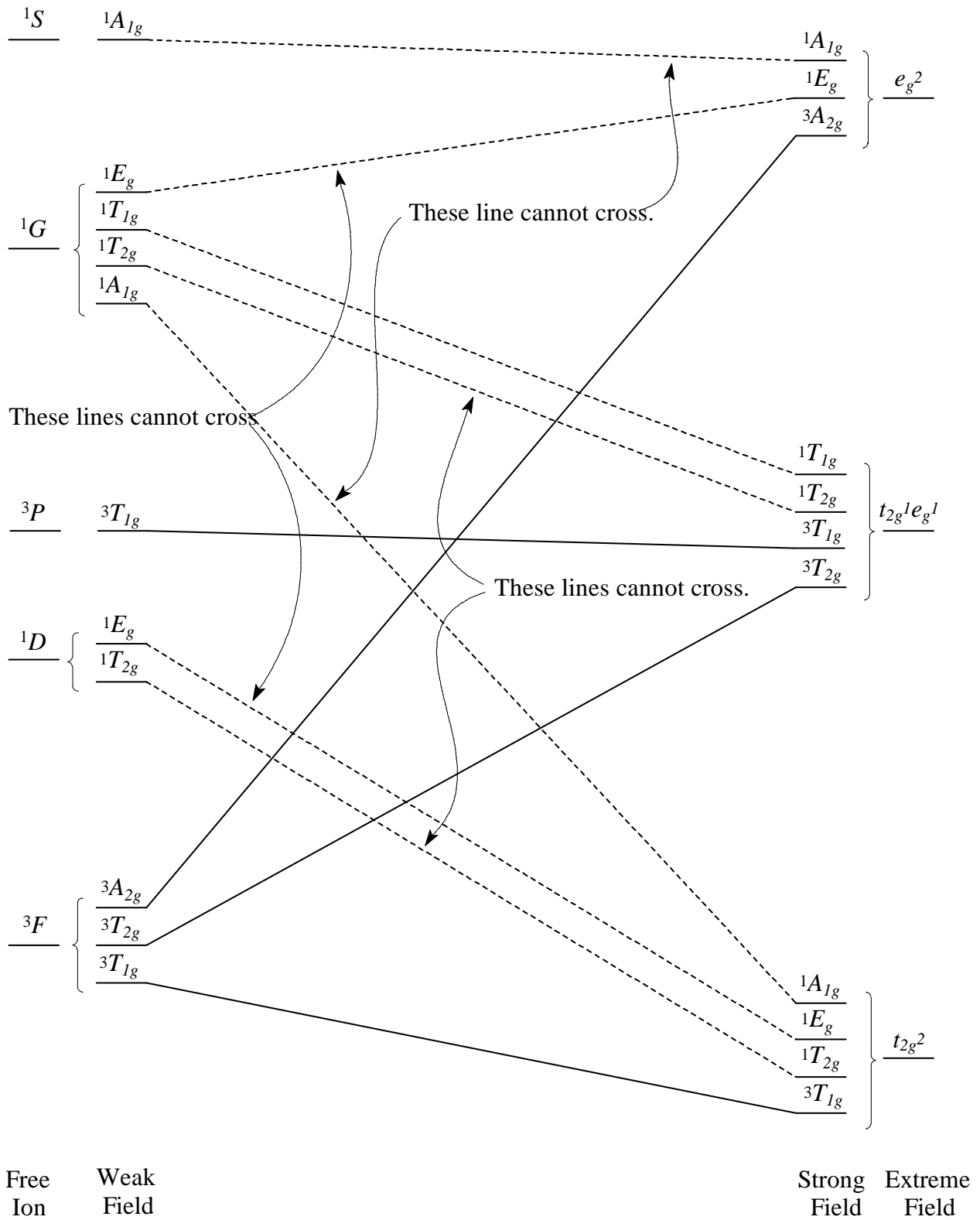
- We now can make the correlation for the remaining two triplet terms on the weak-field side; viz.,  ${}^3T_{1g}(F)$  and  ${}^3T_{1g}(P)$ .
  - From the *noncrossing rule* we conclude that
    - ${}^3T_{1g}(F)$  connects with  ${}^3T_{1g}$  from  $t_{2g}^2$
    - ${}^3T_{1g}(P)$  connects with  ${}^3T_{1g}$  from  $t_{2g}^1 e_g^1$

# Adding the ${}^3T_{1g}$ Correlations

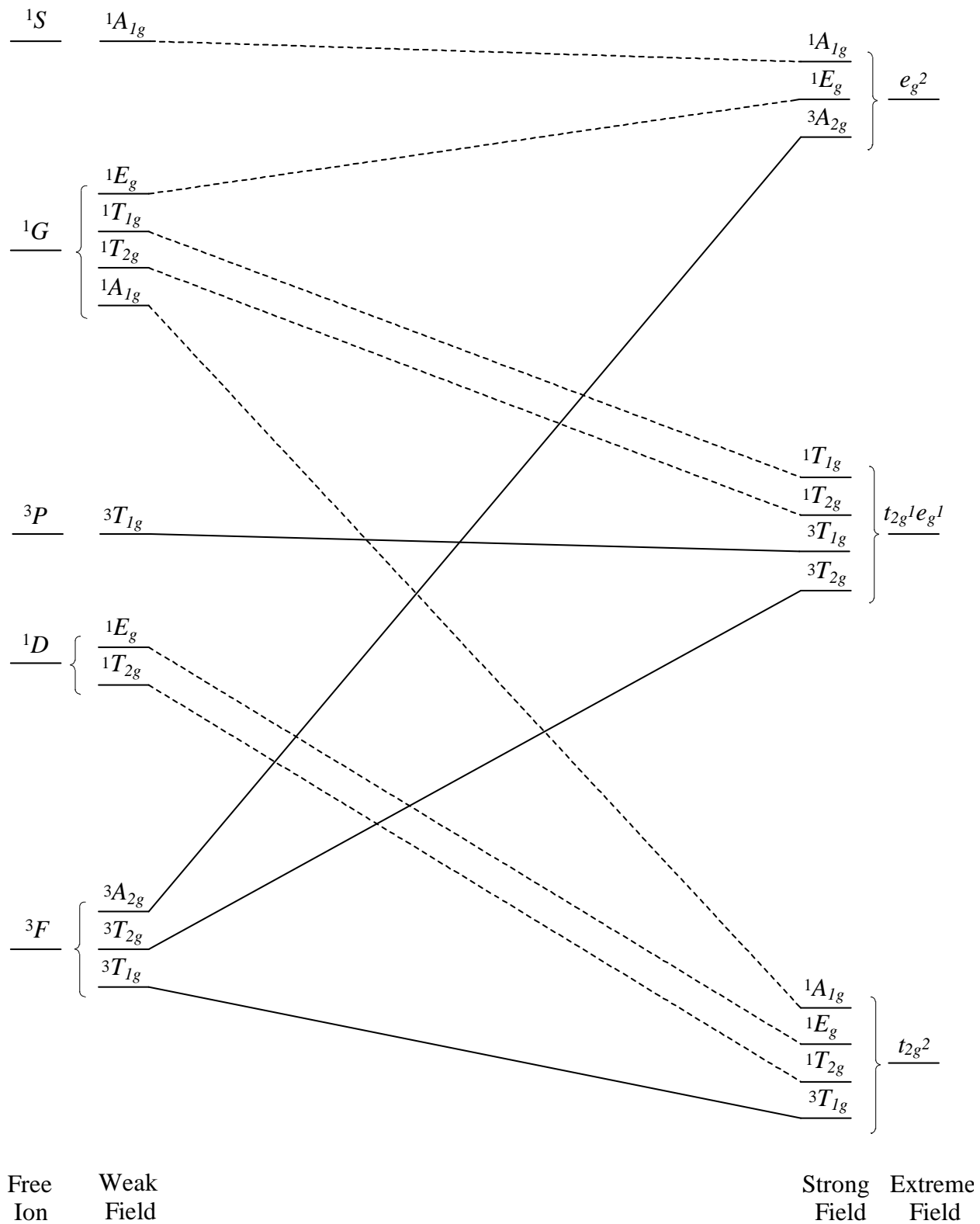


- The remaining correlations between singlet states on both sides can now be made by applying the noncrossing rule.

# Adding the Singlet State Correlations



# The Complete $d^2$ Correlation Diagram



## Other Configurations and Other Fields

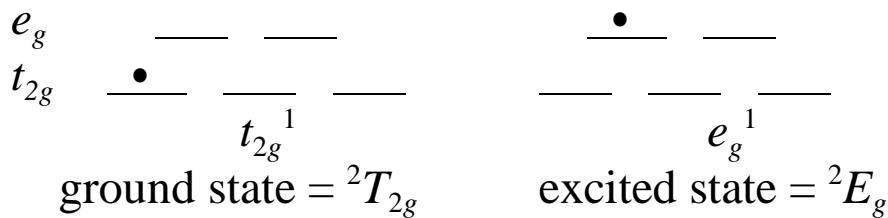
- In principle, we could construct the correlation diagram for any  $d^n$  configuration in an octahedral or other field by taking an approach similar to what we have shown here for the  $d^2$  case.
  - As the number of microstates and terms increases with the number of electrons, the labor of constructing the correlation diagrams “from scratch” becomes considerably more onerous.
  - Fortunately, some general relationships between configurations, terms, and ligand environments minimize the needed effort.

## Ligand Field Splitting for $d^n$ and $d^{10-n}$ Terms

- Free-ion terms for a configuration  $d^n$  and a configuration  $d^{10-n}$  are the same.
  - Each free-ion term splits into a specific collection of terms in any ligand field.
    - ☞ The splitting of terms for a configuration  $d^n$  is identical to that for the configuration  $d^{10-n}$ .
- However, *the order of splitting of a given  $d^n$  term will show the reverse pattern of that of a  $d^{10-n}$  configuration in the same ligand field.*
  - The same ligand-field term that becomes more stable (moves to lower energy) for a  $d^n$  ion will become less stable (moves to higher energy) for a  $d^{10-n}$  ion as  $\Delta_o$  changes.

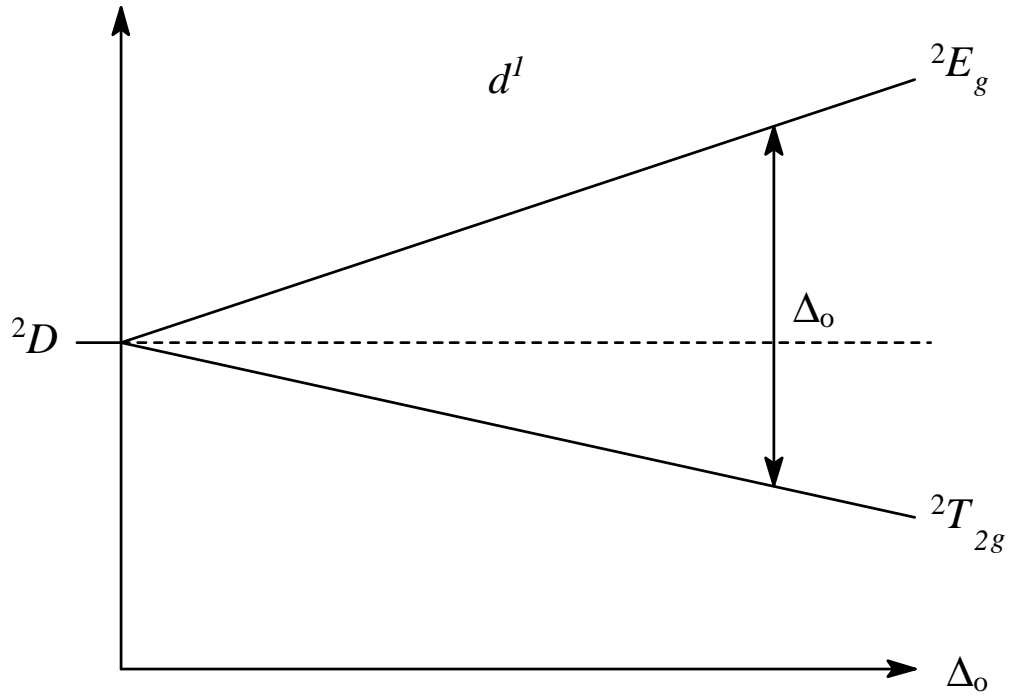
## Term Splitting of $d^1$ and $d^9$ $d^1$ Correlation Diagram

- Both configurations give rise to a  ${}^2D$  free-ion term, which is split into  ${}^2E_g$  and  ${}^2T_{2g}$  terms in an octahedral field.
- In the  $d^1$  case, the two terms in the octahedral field arise from the following configurations



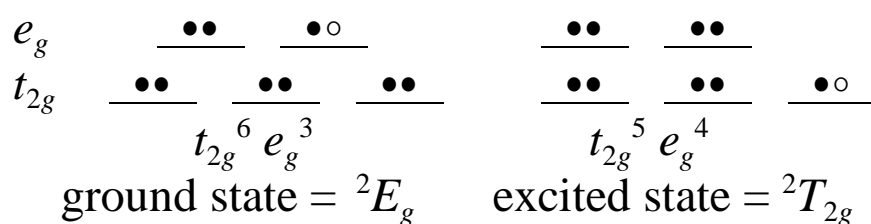
- As the  $\Delta_o$  gap between  $t_{2g}$  and  $e_g$  orbitals increases with increasing field strength the  ${}^2T_{2g}$  term will become more stable and the  ${}^2E_g$  term will become less stable.
  - Thus, the separation between the two states will increase.
  - The separation is numerically equal to  $\Delta_o$ , the magnitude of the field.
  - Relative to the energy of the  ${}^2D$  free-ion term, the  ${}^2T_{2g}$  term will be stabilized by  $-(2/5)\Delta_o$  and the  ${}^2E_g$  term will be destabilized by  $+(3/5)\Delta_o$ .

# $d^1$ Correlation Diagram



## Term Splitting of $d^1$ and $d^9$ $d^9$ Correlation Diagram

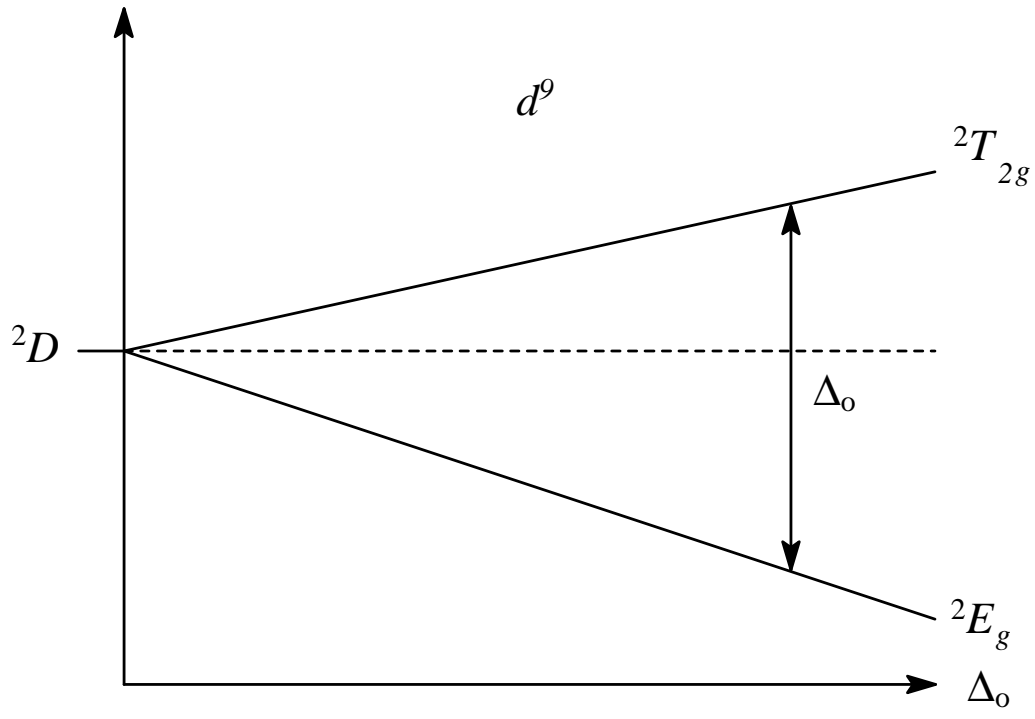
- The same terms will arise from the  ${}^2D$  free-ion term of a  $d^9$  configuration, but they now correspond to the following two configurations in the octahedral field:



(◦ = hole = absence of an electron in an orbital)

- Using the hole as a marker, note that the ground state  $t_{2g}^6 e_g^3$  configuration consists of two equivalent orbital assignments and therefore must correspond to the  ${}^2E_g$  term.
  - Likewise, the three possible orbital assignments for the hole in the  $t_{2g}^5 e_g^4$  configuration show that it gives rise to the  ${}^2T_{2g}$  term.
- ☞ Thus, the energy ordering of the terms for  $d^9$  is the reverse of the  $d^1$  case.
- For the  $d^9$  ( $O_h$ ) case the  ${}^2E_g$  term is stabilized by  $-(3/5)\Delta_o$  and the  ${}^2T_{2g}$  term is destabilized by  $+(2/5)\Delta_o$ .

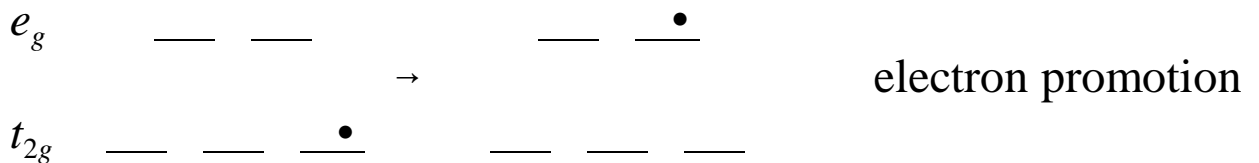
# $d^9$ Correlation Diagram



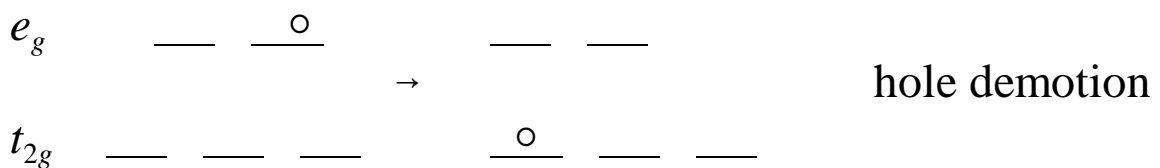
## Hole Formalism

- The relationship between  $d^n$  and  $d^{10-n}$  term splittings in the same symmetry ligand field is sometimes called the *hole formalism*.
  - The name comes from seeing  $d^n$  as a configuration of  $n$  electrons and  $d^{10-n}$  as a configuration of  $n$  positive holes (equivalent to positrons).
  - A configuration of  $n$  electrons will interact with a ligand field in the same way as a configuration of  $n$  positrons, except that repulsions in the former case become attractions in the latter case.

Example: For a  $d^1$  ion in an octahedral field a transition from the ground state to the excited states involves promoting the electron by  $\Delta_o$ .



The same kind of transition for a  $d^9$  ion involves demoting the hole by  $\Delta_o$ .



## Applying Hole Formalism to Correlation Diagrams

- The correlation diagram for  $d^{10-n}$  can be obtained by reversing the order of the sets of terms for the various  $t_{2g}e_g$  configurations on the strong field side of the  $d^n$  diagram, relabeling for the appropriate  $d^{10-n}$  configurations, and redrawing the connecting lines, paying attention to the noncrossing rule.
- We can use this principle with our  $d^2 O_h$  correlation diagram to obtain the  $d^8$  correlation diagram.

- Note that in the  $d^2$  diagram the sets of terms for the three  $t_{2g}e_g$  configurations on the strong-field (right) side are ordered

$$[{}^3T_{1g}, {}^1T_{2g}, {}^1E_g, {}^1A_{1g}] < [{}^3T_{2g}, {}^3T_{1g}, {}^1T_{2g}, {}^1T_{1g}] < [{}^3A_{2g}, {}^1E_g, {}^1A_{1g}]$$

- The order will be reversed for the  $d^8$  diagram:

$$[{}^3A_{2g}, {}^1E_g, {}^1A_{1g}] < [{}^3T_{2g}, {}^3T_{1g}, {}^1T_{2g}, {}^1T_{1g}] < [{}^3T_{1g}, {}^1T_{2g}, {}^1E_g, {}^1A_{1g}]$$

- The ordering of ligand-field terms within each set is arbitrarily chosen for convenience in drawing the diagrams and has no significance regarding relative energies within the set.



## Relationship Between Octahedral and Tetrahedral Fields

- Ligand-field term symbols for the states in a tetrahedral field arising from any  $d^n$  free-ion term are the same as those in an octahedral field, except the labels for the tetrahedral terms omit the subscript  $g$  notation.
- The energies of the new terms in the tetrahedral field have an inverted order.
  - Analogous to the splitting of  $d$  orbitals into  $e$  and  $t_2$  levels in a tetrahedral field, which is the inverse of the splitting into  $t_{2g}$  and  $e_g$  levels in an octahedral field.
  - The tetrahedral and octahedral fields have similar but opposite effects on the  $d$  orbitals.
    - The same is true for the terms arising from  $d^n$  configurations.

## Relationship Between Octahedral and Tetrahedral Correlation Diagrams

- The correlation diagram for  $d^n (T_d)$  can be obtained by reversing the order of the sets of terms for the various  $t_{2g}e_g$  configurations on the strong field side of the  $d^n (O_h)$  diagram, relabeling for the appropriate  $d^n$  tetrahedral configurations, omitting the subscript  $g$  notations from all terms, and redrawing the connecting lines, paying attention to the noncrossing rule.
  - This is essentially the same process we have seen for configurations related by hole formalism.
  - The correlation diagram for  $d^n (T_d)$  is qualitatively the same as that for  $d^{10-n} (O_h)$ , except for minor changes in labels of configurations and term symbols.
- ☞ For example, the  $d^2 (O_h)$  correlation shown is essentially the same as the correlation for  $d^8 (T_d)$ .
- ☞ The  $d^8 (O_h)$  correlation shown is essentially the same as the correlation for  $d^2 (T_d)$ .

## Tanabe-Sugano Diagrams

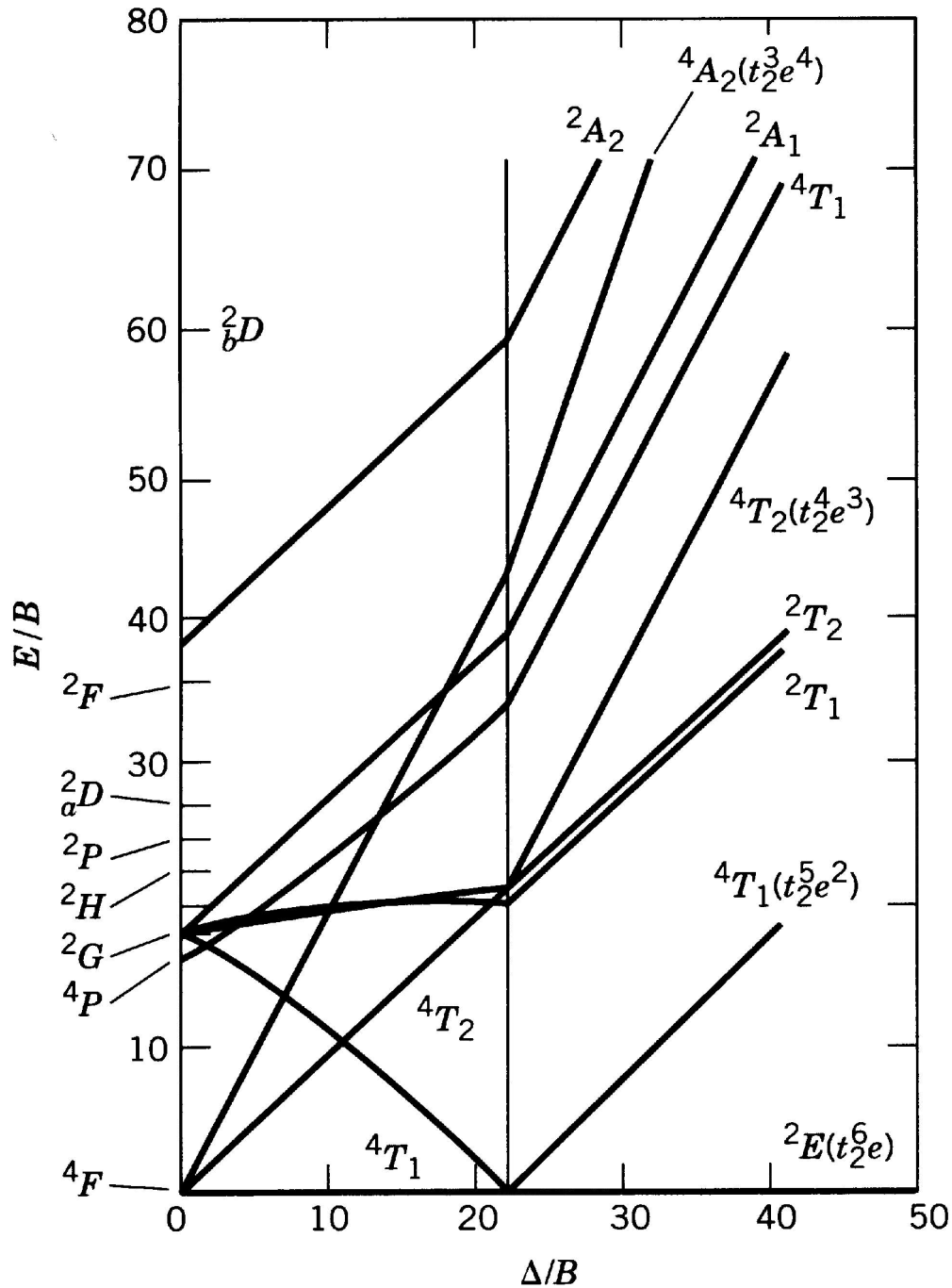
- Most chemists refer to a more detailed set of semi-empirical diagrams for octahedral complexes, originally developed by Yukito Tanabe and Satoru Sugano in 1954.
  - A complete set of these diagrams for octahedral complexes of metal ions with the configurations  $d^2$  through  $d^8$  is shown in Appendix D.
- Tanabe-Sugano diagrams plot term energy versus field strength.
  - Energies of all states are plotted relative to the energy of the ground state term; i.e., the ground state energy forms the abscissa of the plot.
  - Term energies and field strengths are expressed as the variables  $E/B$  and  $\Delta/B$ , respectively, where  $B$  is the *Racah parameter*.
    - The Racah parameter is a measure of the interelectronic repulsion and is used to measure the energy difference between states of the same spin multiplicity.

Example:  $\text{Co}^{2+}$ ,  $3d^7$ , has a difference between the  ${}^4F$  and  ${}^4P$  free-ion terms of  $15B \approx 14,500 \text{ cm}^{-1}$ .

- By using the appropriate values of the Racah parameter, the Tanabe and Sugano diagrams can be used with a variety of metal ions and complexes.

# Tanabe-Sugano Diagram for $d^7$

$d^7$   $B = 970 \text{ cm}^{-1}$  for Co(II)



## Notes on $d^7$ Tanabe-Sugano Diagram

- The  $d^7$  diagram, like all such diagrams for configurations that may be either high spin or low spin, has a perpendicular line near the middle marking the change in spin state.
  - To the left of the line (low field strength, high spin), the ground state is  ${}^4T_1$ , emerging from the free-ion  ${}^4F$  term.
  - To the right of the line (high field strength, low spin), the ground state is  ${}^2E$ , and therefore becomes the abscissa beyond the spin-state crossover point.
  - The high-spin  ${}^2E$  ground state is a continuation of the line for the low-spin excited state  ${}^2E$  term, which emerges from the  ${}^2G$  free-ion term.
  - The line for the former high-spin ground state  ${}^4T_1$  term ascends as an excited state on the low-spin (right) side of the diagram.

## Making Sense of the Confusion

- Sometimes it can be difficult to trace back the free-ion origin of some of the octahedral terms, particularly on the diagrams for high-spin/low-spin configurations.
- Remember: *the spin multiplicities of the split terms must match those of the free-ion terms.*
  - Failure to recognize this has caused some texts to erroneously render the  $d^6$  diagram with the low-spin ground state  $^1A_1$  term originating from the  $^3D$  free-ion term, rather than the correct  $^1I$  term.

