# **Crystal Field Theory History**

- 1929 Hans Bethe Crystal Field Theory (CFT)
  - Developed to interpret color, spectra, magnetism in crystals
- 1932 J. H. Van Vleck CFT of Transition Metal Complexes
  - Champions CFT to interpret properties of transition metal complexes
  - Show unity of CFT, VB, and MO approaches
- 1932 L. Pauling and J. C. Slater VB theory
  - Apply hybrid orbital concepts to interpret properties of transition metal complexes
  - Becomes dominant theory to explain bonding and magnetism until 1950s
  - Can't explain colors and visible spectra
- 1952 L. E. Orgel Revival of CFT and development of Ligand Field Theory (LFT)
  - Slowly replaces VB theory
  - Explains magnetism and spectra better
- 1954 Y. Tanabe and S. Sugano Semi-quantitative term splitting diagrams
  - Used to interpret visible spectra
- 1960s CFT, LFT, and MO Theories
  - Used in conjunction with each other depending on the level of detail required
  - MO used for most sophisticated and quantitative interpretations
  - LFT used for semi-quantitative interpretations
  - CFT used for everyday qualitative interpretations

### **CFT Principles**

- CFT takes an electrostatic approach to the interaction of ligands and metal ions.
  - In purest form it makes no allowances for covalent M–L bonding.
- CFT attempts to describe the effects of the Lewis donor ligands and their electrons on the energies of *d* orbitals of the metal ion.
- We will consider the case of an octahedral  $ML_6(O_h)$  complex first and then extend the approach to other complex geometries.

### Energies of *d* orbitals in an Octahedral Complex

- Consider a spherical field  $(R_3)$  equivalent to six electron pairs surrounding a central metal ion, M.
  - Ligand-metal electron repulsions will perturb the energies of the five degenerate *d* orbitals, making them rise in energy.



- Localize the six ligand pairs into the positions of an octahedron  $(R_3 \rightarrow O_h)$ .
  - Five-fold degeneracy among *d* orbitals will be lifted, in keeping with the direct product listings in the *O<sub>h</sub>* character table.
  - The  $d_{xy}$ ,  $d_{yz}$ ,  $d_{xz}$  orbitals constitute a triply degenerate set of  $t_{2g}$  symmetry.
  - The  $d_{x^2-y^2}$ ,  $d_{z^2}$  orbitals constitute a doubly degenerate set of  $e_g$  symmetry.<sup>1</sup>
- The energies of the  $t_{2g}$  orbitals and  $e_g$  orbitals relative to the perturbed energy of the hypothetical spherical field depend upon their orientation to the six ligand coordination positions.

<sup>&</sup>lt;sup>1</sup>The vector product  $2z^2 - x^2 - y^2$  indicates the *d* orbital more commonly labeled  $z^2$ .

#### d Orbitals in an Octahedral Field

- The  $e_g$  orbitals have lobes that point at the ligands and so will ascend in energy.
- The  $t_{2g}$  orbitals have lobes that lie between ligands and so will descend in energy.











- The energy gap between  $t_{2g}$  and  $e_g$  levels is designated  $\Delta_0$  or 10Dq.
- The energy increase of the  $e_g$  orbitals and the energy decrease of the  $t_{2g}$  orbitals must be balanced relative to the energy of the hypothetical spherical field (sometimes called the *barycenter*).
  - The energy of the  $e_g$  set rises by  $+3/5\Delta_o = +6Dq$  while the energy of the  $t_{2g}$  set falls by  $-2/5\Delta_o = -4Dq$ , resulting in no net energy change for the system.

$$\Delta E = E(e_g)^{\uparrow} + E(t_{2g})^{\downarrow}$$
  
= (2)(+3/5\Delta\_o) + (3)(-2/5\Delta\_o)  
= (2)(+6Dq) + (3)(-4Dq) = 0

- The magnitude of  $\Delta_0$  depends upon both the metal ion and the attaching ligands.
- Magnitudes of  $\Delta_0$  are typically ~100 400 kJ/mol (~8,375 33,500 cm<sup>-1</sup>).<sup>2</sup>

 $<sup>^{2}1 \</sup>text{ kJ/mol} = 83.7 \text{ cm}^{-1}$ 

Electron Filling of  $t_{2g}$  and  $e_g$  Orbitals



- Electrons fill  $t_{2g}$  and  $e_g$  orbitals in an aufbau manner, starting with the  $t_{2g}$  set, in accord with the Pauli exclusion principle and Hund's rule.
  - Spins of successively added electrons are parallel so long as the Pauli exclusion principle allows.
  - At the point when the set of  $t_{2g}$  orbitals is half filled, an additional electron must pair if it is to occupy one of the orbitals of the degenerate set.
  - But if the *mean pairing energy* (*P*) is greater than  $\Delta_0$ , a lower energy state will result by putting the electron in the higher  $e_g$  level.
- For configurations  $d^{1} d^{3}$  and  $d^{8} d^{10}$  there is only one ground state configuration.
- For configurations  $d^4 d^7$  there are two possible filling schemes depending on the magnitudes of *P* and  $\Delta_0$ .
  - A *high spin* configuration minimizes pairing by spreading the electrons across both the  $t_{2g}$  and  $e_g$  levels.
  - A *low spin* configuration minimizes occupying the higher energy  $e_g$  level by pairing electrons in the  $t_{2g}$  level.
- For a given metal ion, the pairing energy is relatively constant, so the spin state depends upon the magnitude of the field strength,  $\Delta_0$ .
  - *Low field* strength ( $\Delta_0 < P$ ) results in a *high-spin* state.
  - *High field* strength ( $\Delta_0 > P$ ) results in a *low-spin* state.

# High- and Low-Spin Configurations for $ML_6 O_h$

$e_{g}$				1		<u>1</u> 1	
$t_{2g}$	 	<u>1</u> 1	<u>1</u> <u>1</u> <u>1</u>	<u>1</u> <u>1</u> 1	1 1 1 L	<u>1</u> <u>1</u> 1	
	$d^{1}$	$d^{2}$	<i>d</i> <sup>3</sup>	d <sup>4</sup> high spin	d <sup>4</sup> low spin	d <sup>5</sup> high spin	d <sup>5</sup> low spin
e <sub>g</sub>	<u>1</u> 1		<u>1</u> 1		<u>1</u> 1	1 1 L	<u> </u>
$t_{2g}$	1 1 1 L		1 1 L 1 L		1 L 1 L 1 L	1 L 1 L 1 L	1   1   1
	d <sup>6</sup> high spin	d <sup>6</sup> low spin	d <sup>7</sup> high spin	d <sup>7</sup> low spin	$d^{8}$	$d^{9}$	$d^{10}$

### **Crystal Field Stabilization Energy (CFSE)**<sup>3</sup>

• Occupancy of electrons in  $t_{2g}$  and  $e_g$  orbitals results in an overall *crystal field stabilization energy* (CFSE), defined for octahedral complexes as<sup>4</sup>

CFSE = 
$$(-0.4n_{t_{2g}} + 0.6n_{e_g})\Delta_0 + pP$$

where  $n_{t_{2g}}$  = number of electrons in  $t_{2g}$  orbitals  $n_{e_g}$  = number of electrons in  $e_g$  orbitals

p = total number of electron pairs

P = mean pairing energy

Free Ion	$O_h$	CFSE Calculation	CFSE
$d^3$	$t_{2g}^{3}$	$(3)(-0.4\Delta_{o})$	-1.2A <sub>o</sub>
$d^{8}$	$t_{2g}^{6}e_{g}^{2}$	$[(6)(-0.4) + (2)(+0.6)]\Delta_0 + 3P$	$-1.2\Delta_{\rm o} + 3P$
$d^7$ low	$t_{2g}^{6}e_{g}^{1}$	$[(6)(-0.4) + (1)(+0.6)]\Delta_0 + 3P$	$-1.8\Delta_{o} + 3P$
$d^7$ high	$t_{2g}^{5}e_{g}^{2}$	$[(5)(-0.4) + (2)(+0.6)]\Delta_0 + 2P$	$-0.8\Delta_{\rm o} + 2P$

- For d<sup>n</sup> cases that could be high- or low-spin, the configuration that results in the lower CFSE for the Δ<sub>o</sub> of the complex is the spin state that is observed.
  - For the hypothetical case  $\Delta_0 = P$ , neither state would be preferred, as the two CFSEs for  $d^7$  illustrate: CFSE $(d^7 \text{ low}) = -1.8\Delta_0 + 3P = -1.8\Delta_0 + 3\Delta_0 = 1.2\Delta_0$

 $CFSE(d^7 \text{ high}) = -0.8\Delta_0 + 2P = -0.8\Delta_0 + 2\Delta_0 = 1.2\Delta_0$ 

• There are no cases for which  $\Delta_0 = P$ .

<sup>&</sup>lt;sup>3</sup>CFSE is also called Ligand Field Stabilization Energy (LFSE).

<sup>&</sup>lt;sup>4</sup>Meissler & Tarr use  $\Pi_c$  for the  $\Delta_o$  term and  $\Pi_e$  for the *P* term in the defining equation. Some sources do not include pairing energy in calculating CFSE.

Complex	$d^n$	$P(\mathrm{cm}^{-1})$	$\Delta_{\rm o}~({\rm cm}^{-1})$	State	CFSE
$[Cr(H_2O)_6]^{2+}$	$d^4$	18,800	13,900	high	$-0.6\Delta_{o}$
$[Fe(H_2O)_6]^{3+}$	$d^5$	24,000	13,700	high	0
$[Fe(H_2O)_6]^{2+}$	$d^6$	14,100	10,400	high	$-0.4\Delta_{\rm o} + P$
$[\operatorname{Fe}(\operatorname{CN})_6]^{4-}$	$d^6$	14,100	33,000	low	$-2.4\Delta_{\rm o} + 3P$
$[CoF_{6}]^{3-}$	$d^6$	16,800	13,000	high	$-0.4\Delta_{\rm o} + P$
$[Co(NH_3)_6]^{3+}$	$d^6$	16,800	23,000	low	$-2.4\Delta_{o}+3P$
$[Co(H_2O)_6]^{2+}$	$d^7$	18,000	9,3000	high	$-0.8\Delta_{\rm o} + 2P$

Values of  $\Delta_0$ , *P*, CFSE and Resulting Spin State

- Values of  $\Delta_0$  depend on both the metal ion and the ligand.
- Most aquo complexes are high spin, because H<sub>2</sub>O is a weak-field ligand.
- Almost all Co<sup>3+</sup> (d<sup>6</sup>) complexes are low spin, including [Co(H<sub>2</sub>O)<sub>6</sub>]<sup>3+</sup>, except [CoF<sub>6</sub>]<sup>3-</sup>, which is high spin.
- Second and third row transition metal ions tend to have low spin states.
  - These ions tend to have larger  $\Delta_0$  values.
  - Larger 4*d* and 5*d* orbitals result in smaller *P* values, owing to lesser electronic repulsions.
  - 4*d* and 5*d* orbitals overlap with ligand orbitals, delocalizing electron density onto the ligands.

For a given metal ion, the magnitude of Δ<sub>o</sub> depends on the ligand and tends to increase according to the following *spectrochemical series*:

$$I^{-} < Br^{-} < Cl^{-} < F^{-} < OH^{-} < C_{2}O_{4}^{2-} < H_{2}O < NH_{3} < en < bipy < phen < CN^{-} \approx CO$$

- *en* = ethylenediamine, *bipy* = 2,2'-bipyradine, *phen* = *o*-phenathroline
- Ligands up through H<sub>2</sub>O are *weak-field ligands* and tend to result in high-spin complexes.
- Ligands beyond H<sub>2</sub>O are *strong-field ligands* and tend to result in low-spin complexes.

## **Tetrahedral Crystal Field Splitting**

- The same considerations of crystal field theory can be applied to  $ML_4$  complexes with  $T_d$  symmetry.
  - In  $T_d$ ,  $d_{xy}$ ,  $d_{yz}$ ,  $d_{xz}$  orbitals have  $t_2$  symmetry and  $d_{x^2-y^2}$ ,  $d_{z^2}$  orbitals have *e* symmetry.
- Relative energies of the two levels are reversed, compared to the octahedral case.
  - $\circ$  No *d* orbitals point directly at ligands.
  - The  $t_2$  orbitals are closer to ligands than are the *e* orbitals. This can be seen by comparing the orientations of the  $d_{x^2-y^2}$  orbital (*e* set) and  $d_{xy}$  orbital ( $t_2$  set) relative to the four ligands.



- The difference results in an energy split between the two levels by  $\Delta_t$  or 10Dq'. Relative to the barycenter defined by the hypothetical spherical field
  - the *e* level is lower by  $-3\Delta_t/5 = -6Dq'$
  - the  $t_2$  level is higher by  $+2\Delta_t / 5 = +4Dq'$

### **Tetrahedral Crystal Field Splitting - cont.**



- In principle, both high and low spin configurations are conceivable for  $d^3-d^6$  ML<sub>4</sub>  $T_d$  complexes.
- With extremely rare exceptions, only high spin configurations are observed.
  - $\circ \Delta_t$  is much smaller than  $\Delta_o$ .
  - For a given ligand at the same M-L distances, it can be shown that  $\Delta_t = (4/9)\Delta_0$ .
  - $\Delta_t \ll P$  in ordinary complexes, so high spin is favored.

• The crystal field stabilization energy for tetrahedral complexes is calculated from the following equation:  $CFSE(T_d) = (-0.6n_e + 0.4n_{t_2})\Delta_t + pP$ 

# **Crystal Field Splitting for Other Geometries**

- We can deduce the CFT splitting of *d* orbitals in virtually any ligand field by
  - Noting the direct product listings in the appropriate character table to determine the ways in which the *d* orbital degeneracies are lifted
  - Carrying out an analysis of the metal-ligand interelectronic repulsions produced by the complex's geometry.
- Sometimes it is useful to begin with either the octahedral or tetrahedral splitting scheme, and then consider the effects that would result by distorting to the new geometry.
  - The results for the perfect and distorted geometries can be correlated through descent in symmetry, using the appropriate correlation tables.
  - Can take this approach with distortions produced by ligand substitution or by intermolecular associations, if descent in symmetry involves a group-subgroup relationship.

### **Crystal Field for Tetragonally Distorted ML<sub>6</sub>**

- A tetragonal distortion to an octahedron results from any change in geometry that preserves a  $C_4$  axis.
  - Tetragonal distortion occurs whenever two *trans* related ligands are differentiated from the remaining four.
- A useful tetragonal distortion to consider involves equally stretching two *trans* related ligands, thereby causing a descent in symmetry  $O_h \rightarrow D_{4h}$ .



- The stretching occurs along the *z* axis, leaving the four positions in the *xy* plane equivalent to each other.
- Ultimately, such a stretching leads to removal of the two ligands, leaving a square planar ML<sub>4</sub> complex.

### Splitting of *d* Orbital Degeneracies – $O_h \rightarrow D_{4h}$

- From a correlation table that links the groups  $O_h$  and  $D_{4h}$  it can be determined that the two  $e_g$  orbitals of the octahedral field become nondegenerate as  $a_{1g}$  and  $b_{1g}$  in the  $D_{4h}$  tetragonal field.
  - From the direct product listings in the  $D_{4h}$  character table

$$a_{1g} = d_{2z^2 - x^2 - y^2} (= d_{z^2})$$
  
$$b_{1g} = d_{x^2 - y^2}$$

- From the correlation table it can also be shown that the degeneracy among the  $t_{2g}$  orbitals in  $O_h$  is partially lifted to become  $b_{2g}$  and  $e_g$  in the  $D_{4h}$  tetragonal field.
  - From the direct product listings in the  $D_{4h}$  character table

$$b_{2g} = d_{xy}$$
$$e_g = (d_{xz}, d_{yz})$$

- Moving the two ligands away from the central metal ion lowers the repulsions between ligand electrons and the metal electrons in those *d* orbitals that have substantial electron distribution along *z*.
  - Thus the energies of the  $d_{xz}$ ,  $d_{yz}$ , and  $d_{z^2}$  orbitals are lowered.
- If we assume that the stretch along *z* is accompanied by a counterbalancing contraction in the *xy* plane, so as to maintain the overall energy of the system, then the orbitals with substantial electron distribution in the *xy* plane will experience increased repulsions.
  - Thus, the  $d_{xy}$  and  $d_{x^2-y^2}$  orbitals rise in energy.

#### **Orbital Splitting from a Stretching Tetragonal Distortion**



- The upper  $e_g$  orbitals of the perfect octahedron split equally by an amount  $\delta_1$ , with the  $d_{x^2-y^2}$  orbital ( $b_{1g}$  in  $D_{4h}$ ) rising by  $+\delta_1/2$  and the  $d_{z^2}$  orbital ( $a_{1g}$  in  $D_{4h}$ ) falling by  $-\delta_1/2$ .
- The lower  $t_{2g}$  orbitals of the perfect octahedron split by an amount  $\delta_2$ , with the  $d_{xy}$  orbital ( $b_{2g}$  in  $D_{4h}$ ) rising by  $+2\delta_2/3$ , and the degenerate  $d_{xz}$  and  $d_{yz}$  orbitals ( $e_g$  in  $D_{4h}$ ) falling by  $-\delta_2/3$ .

#### Magnitudes of the $\delta_1$ and $\delta_2$ Splittings

- Both the  $\delta_1$  and  $\delta_2$  splittings, which are very small compared to  $\Delta_0$ , maintain the barycenters defined by the  $e_g$  and  $t_{2g}$  levels of the undistorted octahedron.
  - The energy gap  $\delta_1$  is larger than that of  $\delta_2$ , because the  $d_{x^2-y^2}$  and  $d_{z^2}$  orbitals are directed at ligands.
  - The distortion has the same effect on the energies of both the  $d_{x^2-y^2}$  and  $d_{xy}$  orbitals; i.e.  $\delta_1/2 = 2\delta_2/3$ .
    - As a result, the energies of both the  $d_{x^2-y^2}$  and  $d_{xy}$  rise in parallel, maintaining a separation equal to  $\Delta_0$  of the undistorted octahedral field.



• Note that  $\delta_1/2 = 2\delta_2/3$  implies that  $\delta_1 = (4/3)\delta_2$ .

#### **Square Planar ML<sub>4</sub> Complexes**

- If we imagine continuing the stretching of M-L bonds along z, the orbital splittings will become progressively greater, producing successively larger values of  $\delta_1$  and  $\delta_2$ .
- Eventually the two ligands will be removed, resulting in a square planar ML<sub>4</sub> complex.
- At some point before this extreme the  $a_{1g}(d_{z^2})$  level may cross and fall below the  $b_{2g}(d_{xy})$  level, resulting in the following splitting scheme.<sup>5</sup>



<sup>&</sup>lt;sup>5</sup>The ordering of the lower four *d* orbitals probably varies among square planar complexes and has been the subject of much debate. See A. B. P. Lever, *Inorganic Electron Spectroscopy*, 2nd ed., Elsevier, Amsterdam, 1984, p. 537ff. and references therein; and, J. Börgl, M. G. Campbell, and T. Ritter, *J. Chem. Educ.*, **2016**, *93*, 118.

### $\mathrm{ML}_{4}\left(D_{4h}\right)$ vs. $\mathrm{ML}_{4}\left(T_{d}\right)$

- Most square planar complexes are  $d^8$  and less often  $d^9$ .
- In virtually all  $d^8$  cases a low spin configuration is observed, leaving the upper  $b_{1g}(d_{x^2-y^2})$  level vacant in the ground state.
  - This is expected, because square planar geometry in firstrow transition metal ions is usually forced by strong field ligands.
  - Strong field ligands produce a large  $\Delta_0$  value.
  - The energy gap between the  $b_{2g}(d_{xy})$  and  $b_{1g}(d_{x^2-y^2})$  levels is equivalent to  $\Delta_0$ .
    - A large  $\Delta_0$  value favors pairing in the  $b_{2g}(d_{xy})$  level, a low-spin diamagnetic configuration for  $d^8$ .
- Tetrahedral  $d^8$  is a high-spin paramagnetic configuration  $e^4t_2^4$ .
  - ML<sub>4</sub>  $(D_{4h})$  and ML<sub>4</sub>  $(T_d)$  can be distinguished by magnetic susceptibility measurements.
- Ni<sup>2+</sup> ion tends to form square planar, diamagnetic complexes with strong-field ligands (e.g., [Ni(CN)<sub>4</sub>]<sup>2-</sup>), but tends to form tetrahedral, paramagnetic complexes with the weaker-field ligands (e.g., [NiCl<sub>4</sub>]<sup>2-</sup>).
- With second and third row transition metal ions the Δ<sub>o</sub> energies are inherently larger, and square planar geometry can occur even with relatively weak field ligands; e.g., square planar [PtCl<sub>4</sub>]<sup>2-</sup>.